Supersymmetry: A status report

Charan Fest

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Congratulations to Charan on his 60th birthday.

Wish much deeper understanding of SUSY SO(10) GUTS from him. Over the last few years, his work on Supersymmetric SO(10) GUTs has been outstanding..perhaps for the first time we seem to have a solid effort towards a "realistic" grand unified theory.

I learnt supersymmetry from Parthasarathy Majumdar and Charan Aulakh on this campus about 18 years ago !

Charan is terrific teacher !!

Charan serves as an inspiration for many of us who have learnt supersymmetry from him.



Supersymmetric Standard Model -1

Jons 40000	00000000
noton - ^^^	mm
V^{\pm} -~~	m
gs- up	www
s-down -	
	Jons \sim Ioton \sim V^{\pm} \sim qs -up \sim is-down \sim

Supersymmetric Standard Model Spectrum -2

Features of SUSY

- It is a technically Natural theory. (Romesh's talk)
- Its calculable and thus in principle, predictable.
- Dark Matter candidate if R-parity is conserved.
- Gauge coupling unification (GUTs with neutrino masses and mixing)
- Lightest Higgs boson can be SM -like in regions of parameter space.



The Higgs bump at LHC CMS $\sqrt{s} = 7$ TeV, L = 5.1 fb⁻¹ $\sqrt{s} = 8$ TeV, L = 5.3 fb⁻¹ Events / 1.5 GeV 0001 Unweighted 120 130 m_{γγ} (GeV) Data S+B Fit **B** Fit Component $\pm 1\sigma$ $\pm 2 \sigma$ 120 130 140 150 110 $m_{\gamma\gamma}$ (GeV)

Speed breakers to Zero Stop mixing ??

Upper bound on Light Higgs (one loop)

 $m_t(m_{SUSY}) \approx 157 \text{ GeV}$

$$m_h^2 = m_Z^2 \cos^2 2\beta + \Delta m_h^2$$

$$\Delta m_h^2 \simeq \frac{3g_2^2 m_t^4}{8\pi^2 M_W^2} \left[\log\left(\frac{m_{\tilde{t}_1} m_{\tilde{t}_2}}{m_t^2}\right) + \frac{X_t^2}{m_{\tilde{t}_1} m_{\tilde{t}_2}} \left(1 - \frac{X_t^2}{12m_{\tilde{t}_1} m_{\tilde{t}_2}}\right) \right]$$

for m_{SUSY} = 1 TeV, we have an upper bound of 135 GeV

pretty robust prediction.

Fixed Order

phenomenological models



Abrey et al. 1112.3028; 2012 updates

For zero mixing, we need multi TeV Stops !!!

Other option is to have maximal mixing : $|X_t| \sim \sqrt{6}M_S$

One loop terms +
dominant 2-loop contribution due to top-stop loops
$$\Pi_{\phi_1}^{(2-\text{loop})}(0) = 0 \qquad \qquad \Pi_{\phi_1\phi_2}^{(2-\text{loop})}(0) = 0$$

$$\Pi_{\phi_2}^{(2-\text{loop})}(0) = \frac{G_F \sqrt{2}}{\pi^2} \frac{\alpha_s}{\pi} \frac{\bar{m}_t^4}{\sin^2 \beta} \left[4 + 3\log^2 \left(\frac{\bar{m}_t^4}{M_S^4} \right) + 2\log \left(\frac{\bar{m}_t^4}{M_S^4} \right) - 6\frac{X_t}{M_S} - \frac{X_t^2}{M_S^2} \left\{ 3\log \left(\frac{\bar{m}_t^2}{M_S^2} \right) + 8 \right\} + \frac{17}{12} \frac{X_t^4}{M_S^4} \right]$$

$$\bar{m}_t = \bar{m}_t(m_t) \approx \frac{m_t^{\text{pole}}}{1 + \frac{4}{3\pi}\alpha_s(m_t)} \qquad + \mathcal{O}(G_F^2 m_t^6)$$
 Heir

Heinemeyer et.al, 9812472

dominant 2-loop correction increases the lightest Higgs mass <10 GeV to the tree-level, assuming the sparticles are <1 TeV (in no-mixing scenario).

Fixed Order

3-loop correction calculated up to $O(\alpha_t \alpha_s^2)$

keeping only the leading terms

no mixing in the stop sector $\Rightarrow X_t = 0$

$$\Delta m_h^{3-\text{loop}} \approx 500 \text{ MeV}$$

 $\sim m_{\star}^4$

Harlander et al. '08 Martin '07

Most Publicly available spectrum generators calculate the CP-even Higgs spectrum at the 2-loop order.

Fixed Order

Theoretical Status of the Higgs mass computation



T.Hahn et. al, arXiv: 1312.4937. Buchmueller et. al, arXiv:1312.5233 Draper et. al 1312.5743

- If LHC discovers light stops (less than TeV) and they are strongly mixed: then MSSM structure is true.
- If LHC discovers light stops and they have zero mixing, it points to structures beyond MSSM (like NMSSM , D-terms etc..)

Charge and Colour breaking Minima

Chowdhury, Godbole, Mohan, Vempati, arXiv: 1310.1932 JHEP

$$\mathcal{V}_{3} = \left(m_{H_{u}}^{2} + \mu^{2}\right)|H_{u}|^{2} + m_{\tilde{t}_{L}}^{2}|\tilde{t}_{L}|^{2} + m_{\tilde{t}_{R}}^{2}|\tilde{t}_{R}|^{2} + (y_{t}A_{t}H_{u}^{*}\tilde{t}_{L}\tilde{t}_{R} + \text{c.c.}) + y_{t}^{2}\left(|\tilde{t}_{L}\tilde{t}_{R}|^{2} + |H_{u}\tilde{t}_{L}|^{2} + |H_{u}\tilde{t}_{R}|^{2}\right) + \frac{g_{1}^{2}}{8}\left(|H_{u}|^{2} + \frac{1}{3}|\tilde{t}_{L}|^{2} - \frac{4}{3}|\tilde{t}_{R}|^{2}\right)^{2} + \frac{g_{2}^{2}}{8}\left(|H_{u}|^{2} - |\tilde{t}_{L}|^{2}\right)^{2} + \frac{g_{3}^{2}}{6}\left(|\tilde{t}_{L}|^{2} - |\tilde{t}_{R}|^{2}\right)^{2}$$

$$= \frac{g_{2}^{2}}{8}\left(|H_{u}|^{2} - |\tilde{t}_{L}|^{2}\right)^{2} + \frac{g_{3}^{2}}{6}\left(|\tilde{t}_{L}|^{2} - |\tilde{t}_{R}|^{2}\right)^{2}$$

$$= \frac{g_{2}^{2}}{6}\left(|H_{u}|^{2} - |\tilde{t}_{L}|^{2}\right)^{2} + \frac{g_{3}^{2}}{6}\left(|\tilde{t}_{L}|^{2} - |\tilde{t}_{R}|^{2}\right)^{2}$$

$$= \frac{g_{3}^{2}}{6}\left(|H_{u}|^{2} - |\tilde{t}_{L}|^{2}\right)^{2} + \frac{g_{4}^{2}}{6}\left(|H_{u}|^{2} - |\tilde{t}_{L}|^{2}\right)^{2} + \frac{g_{4}^{2}}{6}\left(|H_{u}|^{2} - |\tilde{t}_{L}|^{2}\right)^{2} + \frac{g_{4}^{2}}{6}\left(|H_{u}|^{2} - |\tilde{t}_{L}|^{2}\right)^{2} + \frac{g_{4}^{2}}{6}\left(|H_{u}|^{$$

Is the universe in a critical parameter SUSY parameter space?

Stability of MSSM vacuum analysis with four fields, the two Higgs fields and the stop fields (considering they are light)



Higgs productions, decays



Light stops, light staus can significantly modify them...



Carena et. al 1205.5842

Signal strengths can be used to constrain



Limits on Stop masses

Adam Falkowski et. al



Flavour Constraints on third generation

Mostly in sbottom sector

Combined B-physics constraints

$$\frac{(\Delta m_{\tilde{f}}^2)_{23}}{m_{\tilde{f}}^2}\lesssim 10^{-2}~{\rm to}~10^{-1}~{\rm for~about~500~GeV~sbottoms}$$

For larger sbottoms 1 TeV or so, the constraints are much weaker !!

Very similar even in the stau sector

Third generation Summary

Higgs is the strongest constraint pushing the limits close to 1 TeV CCB puts constraints on the mixing of the stops !

Direct constraints are around 300-600 GeV

Flavour violation of O(10)% can be allowed as the third generation masses can reach 1 TeV

Summary of CMS SUSY Results* in SMS framework

SUSY 2013





charginos/neutralinos



Figure 1: Production of pair of chargions (left) and sleptons (right) in the pp collision.



Figure 2: Limits on chargino (left) and slepton (right) pair production

Summary of the data

Gluinos are ruled out up to masses 1-1.25 TeV

First two generations should be greater than 800 GeV -1.25 TeV

(especially if degenerate with the gluino mass)

NO significant constraints on Weakly coupled particles

Flavour Constraints on first two generations

Mostly in squark sector

From \Delta mK

$$rac{(\Delta m_{\widetilde{f}}^2)_{12}}{m_{\widetilde{f}}^2}\lesssim 10^{-3}~{
m to}~10^{-4}~{
m for}$$
 about 500 GeV squarks

For larger squarks 1 TeV or so, the constraints are slightly weaker !!

In slepton sector

From mu to e + gamma

$$rac{(\Delta m_{ ilde{f}}^2)_{12}}{m_{ ilde{f}}^2}\lesssim 10^{-6}~{
m to}~10^{-7}~{
m for}$$
 for about 500 GeV sleptons

Mini Split Supersymmetric Spectrum



Flavour violation can be present in the third generation

neutralino can also be heavy about 1 TeV or so (depending on model)

EDM's could be sensitive to this kind of scenario

Model building is very hard in these scenarios !! Typically single scale susy breaking is very unnatural.

Iyer, Sooryanarayana, Vempati, in preparation

Two scale supersymmetry breaking could be viable

Martin, Nojiri, Bhattacherjee and Mohan, and several others

Degenerate gluino and neutralino mass can escape LHC constraints

Present limits at around 500–600 GeV from mono jet events

A closer look at degenerate MSSM



FIG. 2. The parameter space allowed in $A_t \cdot m_{g-2}$ plane by different constraints for $\tan \beta = 10$ (left panel) and $\tan \beta = 40$ (right panel). The horizontal Grey bands show the δa_{μ} favored values of m_{g-2} at 1σ , 2σ and 3σ (from darker to lighter bands). The red band corresponds to a valid (123-131 GeV) Higgs mass region. The green (yellow) regions are allowed by BR $(B \to X_s \gamma)$ at 1σ (2σ) while the region above the dashed line is allowed by BR $(B_s \to \mu^+ \mu^-)$ at 95% C.L.

Dark Matter interesting in degenerate MSSM (DMSSM)

Chowdhury, Patel, Tata, Vempati, to appear



Spectrum has a splitting between the stops

Dark Matter relic density is very small due to multiple co-annihilations however direct detection can perhaps find a signal !!



Implications on Models

Range we chose $m_0 \in [0, 5]$ TeV $\Delta m_H \in \begin{cases} 0\\ [0, 5] \end{cases}$ for mSUGRA for NUHM1 $m_{1/2} \in [0.1, 2] \text{ TeV}$ $A_0 \in [-3m_0, +3m_0]$ $\operatorname{sgn}(\mu) \in \{-,+\}$



M Raidal et. al arxiv/1112.3647 P. Nath et.al and other groups Baer et.al arXiv: 1112.3017

mSUGRA, tanβ = 10



D. Chowdhury, S. Vempati, et. al

M Raidal et. al arxiv/1112.3647 P. Nath et.al and other groups Baer et.al arXiv: 1112.3017



Dighe et.al arXiv: 1112.3017



Only large negative A -terms are allowed at High scale !!

D. Chowdhury, S. Vempati, et. al,

minimal gauge mediation

Minimal Gauge Mediation

No SUSY flavour violation small number of parameters



SUSY broken spontaneously by X







Two loop diagrams contributing to soft masses







A-terms are essentially zero !!!

the A-terms in the gauge mediation are very small !!

So a 125 GeV Higgs is very difficult unless we have a very heavy stop spectrum (beyond LHC)



FIG. 5. Messenger scale required to produce sufficiently large $|A_t|$ for $m_h = 123$ GeV (left) and $m_h = 125$ GeV (right) through renormalization group evolution.

The change required in the messenger scale is a bit too large : almost up to GUT scale

Ways out for Gauge Mediation

(1) Have Yukawa mediation in addition to gauge mediation. This can be achieved by having matter-messenger fields mixing.

flavour violation !!!!

Delgado, Giudice, Rattazzi et. al, Yanagida et.al

review: Shih et.al, 1303.0228

(2) Have additional matter in the higgs sector.

some amount of Messenger-Matter mixing !

Langacker et. al, Yanagida et. al

(3) Additional strongly coupled sectors

Yanagida et. al

A little more gauge mediation

Say NO to messenger-Matter mixing !!!

Add little more gauge mediation to regular SM gauge group !

Add a singlet !!

(Remember NMSSM does not work in Minimal Gauge Mediation)

 $W = \lambda S H_u H_d + W_{Yuk}$



The RGE generated At is still small !! But the Higgs mass is the in right range !!

Neutrinos can rescue Higgs

Chun, Sooryanarayana, Vempati, to appear

Consider supersymmetric Inverse Seesaw Mechanism



 $m_D \lesssim 0.05 M_R$

From Electroweak precision tests

Perez-Victoria et. al

Neutrinos can rescue Higgs

Chun, Sooryanarayana, Vempati, to appear

Complete 1-loop effective potential corrections including neutrino sector for a general susy breaking Guo et. al sector Shafi et. al



Summary

126 GeV Higgs is compatible with TeV scale MSSM !!! Perhaps it is just around the corner.

But, at the same time, the discovery of Higgs has put severe constraints on known Supersymmetric models even more than direct constraints !!

Whether its mini-split or degenerate or RpV susy we hope to have some idea soon !

Of the models minimal gauge mediation models are the most constrained. But, simple ways can be found without introducing messenger-matter mixing.

For example, we have shown a simple extra U(1) or neutrino couplings can give you the required enhancement without generating large A_t