

Predictions From High Scale Mixing Unification Hypothesis

Rahul Srivastava

The Institute of Mathematical Science
Chennai, India

*Work Done in Collaboration with
G. Rajasekaran, S. Gupta & G. Abbas
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Outline

- 1 Introduction
- 2 High Scale Mixing Unification Hypothesis
- 3 Majorana case
- 4 Dirac Case
- 5 Scale of HSMU and SUSY
- 6 Effect of Phases
- 7 Testing HSMU Hypothesis
- 8 Conclusion and Future Work

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Introduction

- Neutrinos are probably the most mysterious and ill understood of all known particles
- In past neutrinos have thrown up quite a few surprises: They still keep on surprising us !!
- Recent measurements conclusively show $\theta_{13} \neq 0^1$: The latest “surprise”
- Measurement of θ_{13} was long awaited: Provides crucial test of several candidate models
- Is there a “natural” way of understanding non-zero and “relatively large” θ_{13} ?
- In this talk we discuss one such possibility: **High Scale Mixing Unification (HSMU)**

¹T2K, MINOS, DAYA-BAY, RENO and Double Chooz Collaborations

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Current Experimental Scenario

- Global Fits for neutrino oscillation parameters²:

Quantity	Best Fit $\pm 1\text{-}\sigma$	3- σ Range
$\theta_{12}/^\circ$	$33.36^{+0.81}_{-0.78}$	31.09 – 35.89
$\theta_{23}/^\circ$	$40.0^{+2.1}_{-1.5} \oplus 50.4^{+1.3}_{-1.3}$	35.8 – 54.8
$\theta_{13}/^\circ$	$8.66^{+0.44}_{-0.46}$	7.19 – 9.96
$\delta_{\text{CP}}/^\circ$	300^{+66}_{-138}	0 – 360
Δm_{21}^2 (10^{-5} eV ²)	$7.50^{+0.18}_{-0.19}$	7.00 – 8.09
Δm_{31}^2 (10^{-3} eV ²) (N)	$2.473^{+0.070}_{-0.067}$	2.276 – 2.695
Δm_{23}^2 (10^{-3} eV ²) (I)	$2.427^{+0.042}_{-0.065}$	2.242 – 2.649

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Open Questions

- Despite tremendous amount of theoretical and experimental research our understanding of neutrinos is still poor
- Nature of neutrinos: Dirac or Majorana?
- Mass of neutrinos: Hierarchical or quasi degenerate?
- Mass Hierarchy: Normal or Inverted?
- CP violation: δ_{CP} ?
- Octant of θ_{23} : $\theta_{23} < 45^\circ$ or $\theta_{23} > 45^\circ$?
- Why lepton and quark mixing parameters are so different?

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- **Unification of seemingly unrelated phenomenon: An old and quite fruitful notion**
- Has lead to much advancement in our understanding: Electro-Magnetism, Electro-Weak force etc
- Current research: Unification of forces
- Grand Unified Theories (GUTs): Unification of gauge couplings
- Key Ingredient: Quarks and Leptons in same multiplet
- Flavor structure of quarks and leptons: Not totally disconnected
- Interesting possibility: “High Scale” Unification of CKM and PMNS mixing parameters

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Unification of CKM and PMNS mixing parameters

- How is this possible?

- Numerically they are so different from each other!!
- The quark mixing matrix³

$$|U_{\text{CKM}}|_{3\sigma} =$$

$$\begin{pmatrix} 0.97412 \rightarrow 0.97442 & 0.22469 \rightarrow 0.22599 & 0.00337 \rightarrow 0.00366 \\ 0.22455 \rightarrow 0.22585 & 0.97328 \rightarrow 0.97360 & 0.0407 \rightarrow 0.0423 \\ 0.00836 \rightarrow 0.00896 & 0.0399 \rightarrow 0.0415 & 0.999100 \rightarrow 0.999167 \end{pmatrix}$$

- The leptonic mixing matrix⁴

$$|U_{\text{PMNS}}|_{3\sigma} = \begin{pmatrix} 0.795 \rightarrow 0.846 & 0.513 \rightarrow 0.585 & 0.126 \rightarrow 0.178 \\ 0.205 \rightarrow 0.543 & 0.416 \rightarrow 0.730 & 0.579 \rightarrow 0.808 \\ 0.215 \rightarrow 0.548 & 0.409 \rightarrow 0.725 & 0.567 \rightarrow 0.800 \end{pmatrix}$$

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Unification of CKM and PMNS mixing parameters

- Maybe unification at higher scales e.g GUT scale?
- Use RG equations to obtain values at low scale (M_Z)
- Hierarchical nature of quark masses: Quark mixing angles don't change much (SM/MSSM RG running)

- What about neutrino mixing angles?

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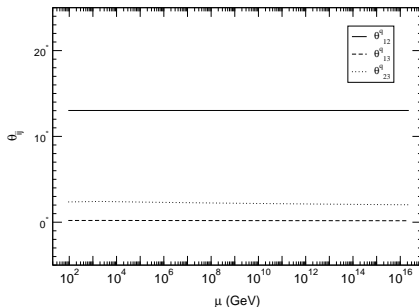
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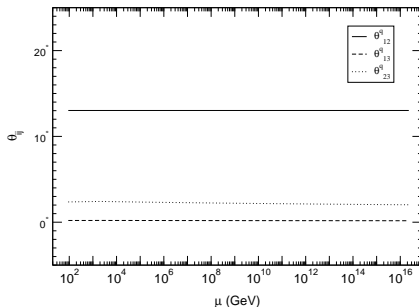
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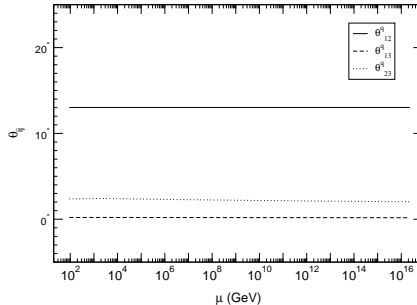
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Radiative Magnification

- High Scale Mixing Unification (HSMU): CKM angles = PMNS angles
- More specifically: For unification at some “High Scale”, say GUT

$$\theta_{12}^{0,q} = \theta_{12}^0 = 13.02^\circ, \quad \theta_{13}^{0,q} = \theta_{13}^0 = 0.17^\circ, \quad \theta_{23}^{0,q} = \theta_{23}^0 = 2.03^\circ$$

- Large radiative magnification of PMNS angles is required

$$\theta_{12}^0 = 13.02^\circ \rightarrow \theta_{12} = 33.36^\circ, \quad \theta_{13}^0 = 0.17^\circ \rightarrow \theta_{13} = 8.66^\circ, \\ \theta_{23}^0 = 2.03^\circ \rightarrow \theta_{23} = 40.0 \oplus 50.4^\circ$$

- Not possible within Standard Model (SM)
- Can be realized within Minimal Supersymmetric Standard Model (MSSM)⁵

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HSMU: General Framework

- **Model independent approach: Assume HSMU at some “High Scale”**
- Details of the “High Scale” theory not needed
- Below High Scale: MSSM + Type-I seesaw mechanism

$$\mathcal{L}_{MSSM+SSI} = \mathcal{L}_{MSSM} + (Y_\nu)_{ij} \nu^{Ci} \mathbf{h}_a^{(u)} \varepsilon^{ab} \mathbf{l}_b^j \Big|_{\theta\theta} + \frac{1}{2} M_{ij} \nu^{Ci} \nu^{Cj} \Big|_{\theta\theta} + h.c.$$

- Effective left handed neutrino mass matrix

$$m_\nu(\mu) = -\frac{v^2}{2} Y_\nu^T(\mu) M^{-1}(\mu) Y_\nu(\mu)$$

- Right handed neutrinos integrated out below their mass threshold

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$$m_\nu(\mu) = -\frac{v^2}{2} Y_\nu^T(\mu) M^{-1}(\mu) Y_\nu(\mu)$$

- Right handed neutrinos integrated out below their mass threshold

HSMU: General Framework

- Model independent approach: Assume HSMU at some “High Scale”
- Details of the “High Scale” theory not needed
- Below High Scale: MSSM + Type-I seesaw mechanism

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HSMU: General Framework

- Below seesaw scale: Effective dimension five neutrino mass operator

$$\mathcal{L}_{MSSM+\kappa} = \mathcal{L}_{MSSM} - \frac{1}{4} \kappa_{ij} l_a^i \epsilon^{ab} \mathbf{h}_b^{(u)} \nu_c^j \epsilon^{cd} \mathbf{h}_d^{(u)} \Big|_{\theta\theta}$$

- Testing HSMU: Need to run down the masses and mixing parameters from High Scale to low scale (M_Z)
- RG running between High Scale and seesaw scale: Using standard MSSM RG equations within framework of Type-I seesaw mechanism
- Below seesaw scale: RG running with dim-5 operator added to MSSM
- Below SUSY breaking scale: RG running with dim-5 operator added to SM

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RG equations: Mass Matrix

- The RG equation of the effective mass operator⁶

$$16\pi^2 \frac{d\kappa}{dt} = C(Y_e^\dagger Y_e)^T \kappa + C\kappa(Y_e^\dagger Y_e) + \alpha\kappa$$

where $t = \ln(\mu/\mu_0)$, μ is the renormalization scale and $C = 1(\frac{-3}{2})$ in MSSM(SM).

- In MSSM and SM α reads

$$\begin{aligned}\alpha_{\text{MSSM}} &= -\frac{6}{5}g_1^2 - 6g_2^2 + 6(y_t^2 + y_c^2 + y_u^2) \\ \alpha_{\text{SM}} &= -3g_2^2 + 2(y_\tau^2 + y_\mu^2 + y_e^2) + 6(y_t^2 + y_b^2 + y_c^2 \\ &\quad + y_s^2 + y_d^2 + y_u^2) + \lambda\end{aligned}$$

where y_f , ($f = \{e, d, u\}$) are the Yukawa couplings, g_i are gauge couplings and λ is SM Higg's quartic coupling.

⁶S. Antusch, J. Kersten, M. Lindner and M. Ratz, Nucl. Phys. B 674, 401 (2003) [hep-ph/0305273]

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Mass Basis

- Parameters of interest: Masses, mixing angles and physical phases

- Need to go to mass basis: $\text{diag}(m_1, m_2, m_3)$

- Parameterization of PMNS matrix:

$$\begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix} \begin{pmatrix} e^{\frac{-i\phi_1}{2}} & 0 & 0 \\ 0 & e^{\frac{-i\phi_2}{2}} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

where $c_{ij} = \cos \theta_{ij}$ and $s_{ij} = \sin \theta_{ij}$ ($i, j = 1, 2, 3$)

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RG equations: Masses

- RG running of masses⁷

$$16\pi^2 \frac{dm_1}{dt} = [\alpha + Cy_\tau^2 (2s_{12}^2 s_{23}^2 + F_1)] m_1 ,$$

$$16\pi^2 \frac{dm_2}{dt} = [\alpha + Cy_\tau^2 (2c_{12}^2 s_{23}^2 + F_2)] m_2 ,$$

$$16\pi^2 \frac{dm_3}{dt} = [\alpha + 2Cy_\tau^2 c_{13}^2 c_{23}^2] m_3 ,$$

- Where F_1 and F_2 are:

$$F_1 = -s_{13} \sin 2\theta_{12} \sin 2\theta_{23} \cos \delta + 2s_{13}^2 c_{12}^2 c_{23}^2 ,$$

$$F_2 = s_{13} \sin 2\theta_{12} \sin 2\theta_{23} \cos \delta + 2s_{13}^2 s_{12}^2 c_{23}^2 .$$

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RG equations: Mass Square Differences

- Easily translated into RGEs for the mass squared differences

$$8\pi^2 \frac{d}{dt} \Delta m_{\text{sol}}^2 = \alpha \Delta m_{\text{sol}}^2 + Cy_\tau^2 [2s_{23}^2 (m_2^2 c_{12}^2 - m_1^2 s_{12}^2) + F_{\text{sol}}] ,$$

$$8\pi^2 \frac{d}{dt} \Delta m_{\text{atm}}^2 = \alpha \Delta m_{\text{atm}}^2 + Cy_\tau^2 [2m_3^2 c_{13}^2 c_{23}^2 - 2m_2^2 c_{12}^2 s_{23}^2 + F_{\text{atm}}] ,$$

- Where

$$F_{\text{sol}} = (m_1^2 + m_2^2) s_{13} \sin 2\theta_{12} \sin 2\theta_{23} \cos \delta \\ + 2s_{13}^2 c_{23}^2 (m_2^2 s_{12}^2 - m_1^2 c_{12}^2) ,$$

$$F_{\text{atm}} = -m_2^2 s_{13} \sin 2\theta_{12} \sin 2\theta_{23} \cos \delta - 2m_2^2 s_{13}^2 s_{12}^2 c_{23}^2 .$$

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RG equations: Angles

- RG running of angles⁸

$$\dot{\theta}_{12} = -\frac{Cy_\tau^2}{32\pi^2} \sin 2\theta_{12} s_{23}^2 \frac{|m_1 e^{i\phi_1} + m_2 e^{i\phi_2}|^2}{\Delta m_{\text{sol}}^2} + \mathbf{O}(\theta_{13})$$

$$\begin{aligned} \dot{\theta}_{13} = & \frac{Cy_\tau^2}{32\pi^2} \sin 2\theta_{12} \sin 2\theta_{23} \frac{m_3}{\Delta m_{\text{atm}}^2 (1 + \zeta)} \times \\ & \times [m_1 \cos(\phi_1 - \delta) - (1 + \zeta) m_2 \cos(\phi_2 - \delta) - \zeta m_3 \cos \delta] + \mathbf{O}(\theta_{13}) \end{aligned}$$

$$\begin{aligned} \dot{\theta}_{23} = & -\frac{Cy_\tau^2}{32\pi^2} \sin 2\theta_{23} \frac{1}{\Delta m_{\text{atm}}^2} [c_{12}^2 |m_2 e^{i\phi_2} + m_3|^2 \\ & + s_{12}^2 \frac{|m_1 e^{i\phi_1} + m_3|^2}{1 + \zeta}] + \mathbf{O}(\theta_{13}) \end{aligned}$$

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RG equations: Dirac Phase

- RG running of Dirac phase⁹

$$\dot{\delta} = \frac{C y_{\tau}^2}{32\pi^2} \frac{\delta^{(-1)}}{\theta_{13}} + \frac{C y_{\tau}^2}{8\pi^2} \delta^{(0)} + \mathbf{O}(\theta_{13}) ,$$

where

$$\begin{aligned} \delta^{(-1)} &= \sin 2\theta_{12} \sin 2\theta_{23} \frac{m_3}{\Delta m_{\text{atm}}^2 (1 + \zeta)} \times \\ &\quad \times [m_1 \sin(\phi_1 - \delta) - (1 + \zeta) m_2 \sin(\phi_2 - \delta) + \zeta m_3 \sin \delta] , \\ \delta^{(0)} &= \frac{m_1 m_2 s_{23}^2 \sin(\phi_1 - \phi_2)}{\Delta m_{\text{sol}}^2} \\ &\quad + m_3 s_{12}^2 \left[\frac{m_1 \cos 2\theta_{23} \sin \phi_1}{\Delta m_{\text{atm}}^2 (1 + \zeta)} + \frac{m_2 c_{23}^2 \sin(2\delta - \phi_2)}{\Delta m_{\text{atm}}^2} \right] \\ &\quad + m_3 c_{12}^2 \left[\frac{m_1 c_{23}^2 \sin(2\delta - \phi_1)}{\Delta m_{\text{atm}}^2 (1 + \zeta)} + \frac{m_2 \cos 2\theta_{23} \sin \phi_2}{\Delta m_{\text{atm}}^2} \right] . \end{aligned}$$

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RG equations: Majorana Phases

- RG running of the physical Majorana phases¹⁰

$$\begin{aligned}\dot{\phi}_1 &= \frac{Cy_\tau^2}{4\pi^2} \left\{ m_3 \cos 2\theta_{23} \frac{m_1 s_{12}^2 \sin \phi_1 + (1 + \zeta) m_2 c_{12}^2 \sin \phi_2}{\Delta m_{\text{atm}}^2 (1 + \zeta)} \right. \\ &\quad \left. + \frac{m_1 m_2 c_{12}^2 s_{23}^2 \sin(\phi_1 - \phi_2)}{\Delta m_{\text{sol}}^2} \right\} + \mathbf{O}(\theta_{13}) , \\ \dot{\phi}_2 &= \frac{Cy_\tau^2}{4\pi^2} \left\{ m_3 \cos 2\theta_{23} \frac{m_1 s_{12}^2 \sin \phi_1 + (1 + \zeta) m_2 c_{12}^2 \sin \phi_2}{\Delta m_{\text{atm}}^2 (1 + \zeta)} \right. \\ &\quad \left. + \frac{m_1 m_2 s_{12}^2 s_{23}^2 \sin(\phi_1 - \phi_2)}{\Delta m_{\text{sol}}^2} \right\} + \mathbf{O}(\theta_{13}) .\end{aligned}$$

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HSMU: Assumptions and Initial Conditions

- **Natural “High Scale”: Grand Unified Theory (GUT) Scale**
- Assume HSMU realized at GUT scale i.e. 2×10^{16} GeV
- Sensitivity to choice of high scale: Discussed in later part of talk
- Choice of seesaw scale: HSMU realized for varied range of seesaw scale
- For sake of definiteness: Choose typical Seesaw scale of 10^{12} GeV
- SUSY breaking scale: 5 TeV
- Dependence on choice of SUSY breaking scale: Discussed in later part of talk
- Larger values of $\tan \beta$: Enhanced magnification
- We choose $\tan \beta = 55$: Dependence on $\tan \beta$ discussed in later part of talk

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- Inverted Hierarchy: No magnification of θ_{23}
- Assume no CP violation in leptonic sector: Dirac as well as Majorana phases taken zero
- CP violating scenario: Effect of phases on HSMU discussed in later part of talk

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Implementing HSMU: Two step process

• Bottom - Up

- Start from known values of gauge couplings, quark mixing angles, masses of quarks and charged leptons at low energies (M_Z)
- Use RG equations: Obtain the corresponding values at high energies
- HSMU hypothesis: Take neutrino mixing angles same as the quark mixing angles at the unification scale

• Top - Down

- Neutrino masses at high scale: Unknown parameters
- Determine these three parameters such that: Low energy values of the oscillation parameters i.e. $\Delta m_{12}^2, \Delta m_{23}^2, \theta_{12}, \theta_{23}$ and θ_{13} agree with their present experimental ranges

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- Start from known values of gauge couplings, quark mixing angles, masses of quarks and charged leptons at low energies (M_Z)
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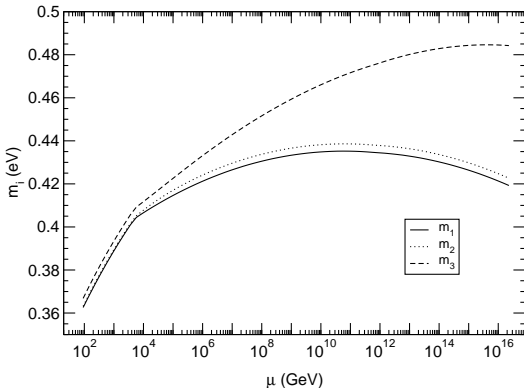
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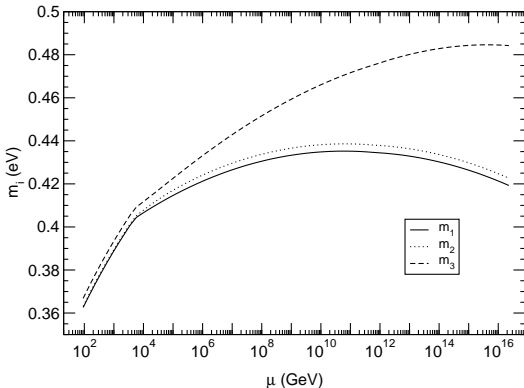
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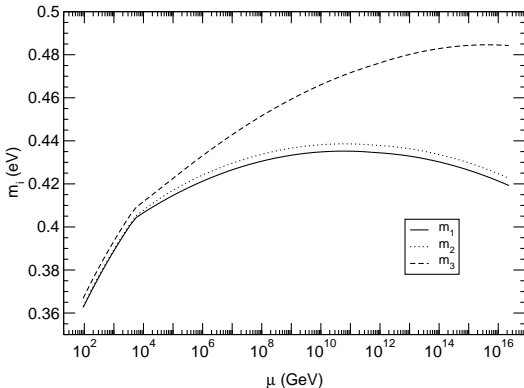
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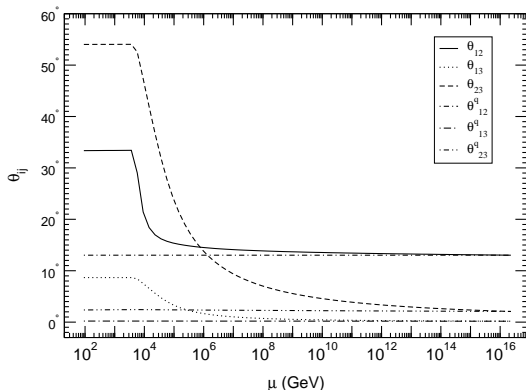
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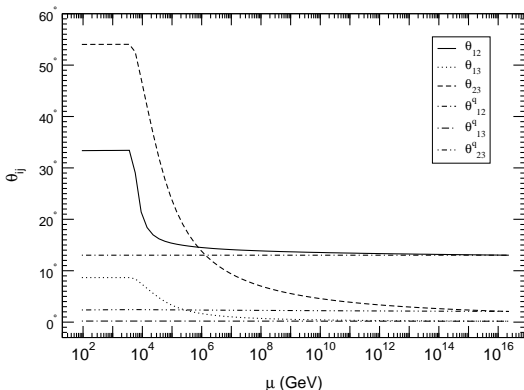


- Hierarchical quark masses:
RG running in quark sector
is almost negligible
- RG running of leptonic mixing
angles

$$\frac{d\theta_{12}}{dt} \propto \frac{m^2}{\Delta m_{21}^2}$$

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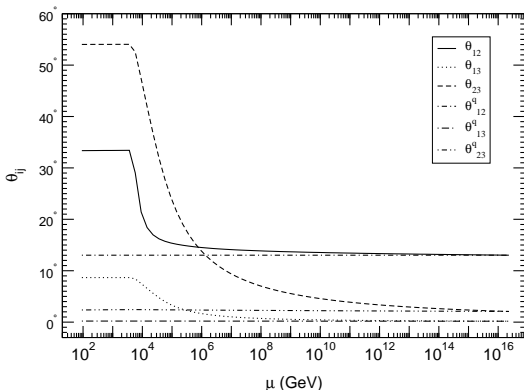
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Numerical results on the evolution of masses and mixing



	I	II	III
$m_1^0(\text{eV})$	0.4196	0.4146	0.4286
$m_2^0(\text{eV})$	0.4230	0.4180	0.4320
$m_3^0(\text{eV})$	0.4843	0.4786	0.4946
$m_1(\text{eV})$	0.3626	0.3582	0.3703
$m_2(\text{eV})$	0.3632	0.3589	0.3709
$m_3(\text{eV})$	0.3668	0.3625	0.3746
$\Delta m_{21}^2(\text{eV}^2)_{RG}$	4.30×10^{-4}	4.49×10^{-4}	4.20×10^{-4}
$\Delta m_{32}^2(\text{eV}^2)_{RG}$	2.67×10^{-3}	2.62×10^{-3}	2.78×10^{-3}
$M_{\tilde{e}}/M_{\tilde{\mu}, \tilde{\tau}}$	1.94	1.84	2.16
$\Delta m_{21}^2(\text{eV}^2)_{th}$	-3.55×10^{-4}	-3.73×10^{-4}	-3.44×10^{-4}
$\Delta m_{32}^2(\text{eV}^2)_{th}$	-2.74×10^{-4}	-2.16×10^{-4}	-3.82×10^{-4}
$\Delta m_{21}^2(\text{eV}^2)_{tot}$	7.50×10^{-5}	7.58×10^{-5}	7.59×10^{-5}
$\Delta m_{32}^2(\text{eV}^2)_{tot}$	2.40×10^{-3}	2.40×10^{-3}	2.40×10^{-3}
$\theta_{23}/^\circ$	54.03	53.93	54.18
$\theta_{13}/^\circ$	8.66	8.67	8.67
$\theta_{12}/^\circ$	33.36	31.14	35.87

Testing HSMU: Values of $m_{\beta\beta}$ and m_β

- Threshold corrections needed, to obtain Δm_{21}^2 within $3\text{-}\sigma$ range¹¹.
- The various entries in the table also highlight the correlations between low scale neutrino oscillation parameters.
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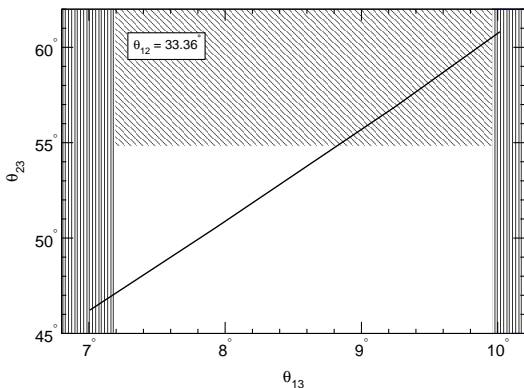
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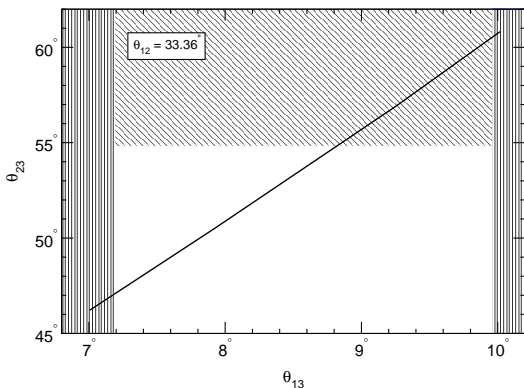
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Octant of θ_{23}



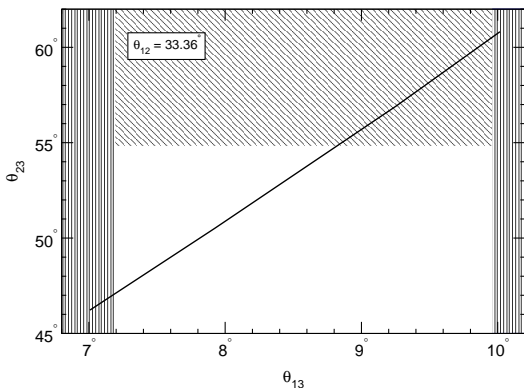
- Since $\frac{d\theta_{13}}{dt}, \frac{d\theta_{23}}{dt} \propto \frac{m^2}{\Delta m_{32}^2}$:
RG evolution of θ_{13} and θ_{23} correlated
- All other oscillation parameters are at their best-fit values.
- Non maximal θ_{23} i.e. $\theta_{23} > 45^\circ$:
Lies in second octant for the whole 3- σ range of θ_{13} .
- Shaded regions lie outside 3- σ range.
- For a fixed value of θ_{13} :
Effect of variation of θ_{12} on θ_{23} is negligible.

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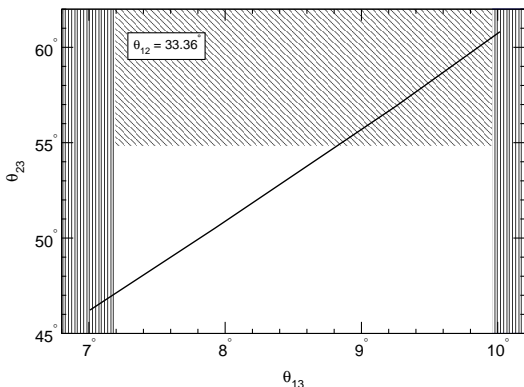
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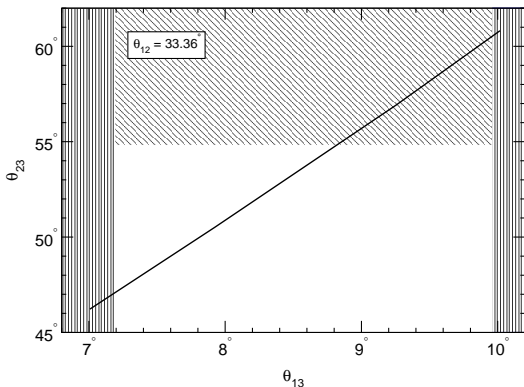
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- Several important predictions:
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Outline

- 1 Introduction
- 2 High Scale Mixing Unification Hypothesis
- 3 Majorana case
- 4 Dirac Case**
- 5 Scale of HSMU and SUSY
- 6 Effect of Phases
- 7 Testing HSMU Hypothesis
- 8 Conclusion and Future Work

Nature of Neutrinos: Dirac or Majorana?

- One of the most important open questions in neutrino physics: Whether neutrinos are Dirac or Majorana particles
- Answering this question: Essential to find the underlying theory of neutrino masses and mixing.
- Current understanding: Dirac neutrinos as plausible as Majorana ones
- Neutrinoless double beta decay experiments: Dedicated ongoing experiments to determine the nature of neutrinos.
- No conclusive evidence: Neutrinoless double beta decay experiments have not seen any signal so far¹².
- Instructive to see if HSMU can be implemented for Dirac Neutrinos as well

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- Answering this question: Essential to find the underlying theory of neutrino masses and mixing.
- Current understanding: Dirac neutrinos as plausible as Majorana ones
- **Neutrinoless double beta decay experiments: Dedicated ongoing experiments to determine the nature of neutrinos.**
- No conclusive evidence: Neutrinoless double beta decay experiments have not seen any signal so far¹².
- Instructive to see if HSMU can be implemented for Dirac Neutrinos as well

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HSMU for Dirac Neutrinos

- **HSMU hypothesis: More natural for Dirac neutrinos than Majorana neutrinos**
- If neutrinos are Majorana particles: The PMNS-matrix has 6 independent parameters; 3-mixing angles, 1-Dirac phase and 2-Majorana phases
- On the other hand CKM-matrix has only 4 independent parameters: 3-mixing angles and 1-Dirac phase
- Clear mismatch between number of parameters on two sides and hence a one-to-one correspondence is impossible
- HSMU for Majorana Case: One has to treat the Majorana phases as free parameters
- Majorana phases influence RG evolution of mixing angles: Predictions subject to choice of Majorana phases
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RG equation: Masses

- RG running of masses¹³

$$16\pi^2 \dot{m}_1 = \left\{ C y_\tau^2 \left[\cos^2 \theta_{12} \cos^2 \theta_{23} \sin^2 \theta_{13} + \sin^2 \theta_{12} \sin^2 \theta_{23} - \frac{1}{2} \cos \delta \sin \theta_{13} \sin(2\theta_{12}) \sin(2\theta_{23}) \right] + \alpha \right\} m_1 ,$$

$$16\pi^2 \dot{m}_2 = \left\{ C y_\tau^2 \left[\sin^2 \theta_{12} \cos^2 \theta_{23} \sin^2 \theta_{13} + \cos^2 \theta_{12} \sin^2 \theta_{23} + \frac{1}{2} \cos \delta \sin \theta_{13} \sin(2\theta_{12}) \sin(2\theta_{23}) \right] + \alpha \right\} m_2 ,$$

$$16\pi^2 \dot{m}_3 = \left\{ C y_\tau^2 \cos^2 \theta_{13} \cos^2 \theta_{23} + \alpha \right\} m_3 .$$

¹³M. Lindner, M. Ratz and M. A. Schmidt, JHEP 0509, 081 (2005), hep-ph/0506280

RG equation: Angles

- RG running of mixing angles¹⁴

$$\begin{aligned}\dot{\theta}_{12} &= \frac{-C y_\tau^2}{32 \pi^2} \frac{m_1^2 + m_2^2}{m_2^2 - m_1^2} \sin(2 \theta_{12}) \sin^2 \theta_{23} + \mathbf{O}(\theta_{13}) , \\ \dot{\theta}_{13} &= \frac{-C y_\tau^2}{32 \pi^2} \frac{1}{(m_3^2 - m_1^2)(m_3^2 - m_2^2)} \{ (m_2^2 - m_1^2) m_3^2 \\ &\quad \cos \delta \cos \theta_{13} \sin(2 \theta_{12}) \sin(2 \theta_{23}) + [m_3^4 - (m_2^2 - m_1^2) \\ &\quad m_3^2 \cos(2 \theta_{12}) - m_1^2 m_2^2] \cos^2 \theta_{23} \sin(2 \theta_{13}) \} , \\ \dot{\theta}_{23} &= \frac{-C y_\tau^2}{32 \pi^2} \frac{[m_3^4 - m_1^2 m_2^2 + (m_2^2 - m_1^2) m_3^2 \cos(2 \theta_{12})]}{(m_3^2 - m_1^2)(m_3^2 - m_2^2)} \sin(2 \theta_{23}) \\ &\quad + \mathbf{O}(\theta_{13}) ,\end{aligned}$$

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RG equation: Dirac Phase

- RG running of Dirac phase¹⁵

$$\dot{\delta} = \dot{\delta}^{(-1)}\theta_{13}^{-1} + \dot{\delta}^{(0)} + \dot{\delta}^{(1)} + \mathbf{O}(\theta_{13}^2),$$

where

$$\dot{\delta}^{(-1)} = \frac{C y_\tau^2}{32 \pi^2} \frac{(m_2^2 - m_1^2) m_3^2}{(m_3^2 - m_1^2)(m_3^2 - m_2^2)} \sin(\delta) \sin(2\theta_{12}) \sin(2\theta_{23}),$$

$$\dot{\delta}^{(0)} = 0$$

$$\dot{\delta}^{(1)} = \frac{C y_\tau^2}{16 \pi^2} \frac{m_2^2 (m_3^2 - m_1^2)^2}{(m_2^2 - m_1^2)(m_3^2 - m_1^2)(m_3^2 - m_2^2)} \cot(\theta_{12}) \sin(2\theta_{23}) \sin \delta + \dots,$$

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Implementing HSMU

- **Same as before**
- Choose: Unification scale = 2×10^{16} GeV, SUSY breaking scale = 5 TeV and $\tan \beta = 55$
- **Bottom - Up**
- Start from known values of gauge couplings, quark mixing angles, masses of quarks and charged leptons at low energies (M_Z)
- Use RG equations: Obtain the corresponding values at high energies
- HSMU hypothesis: Take neutrino mixing angles and phase same as the quark mixing angles and phase at the unification scale
- **Top - Down**
- Neutrino masses at high scale: Unknown parameters
- Determine these three parameters such that: Low energy values of the oscillation parameters i.e. $\Delta m_{12}^2, \Delta m_{23}^2, \theta_{12}, \theta_{23}$ and θ_{13} agrees with their present experimental ranges

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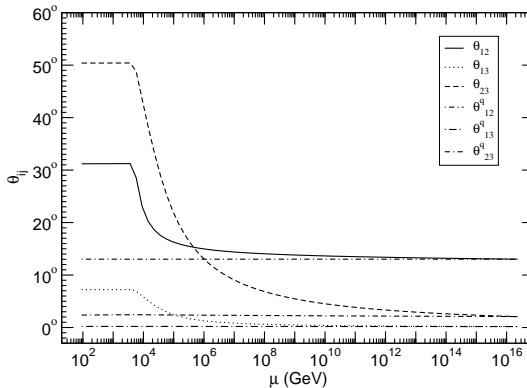
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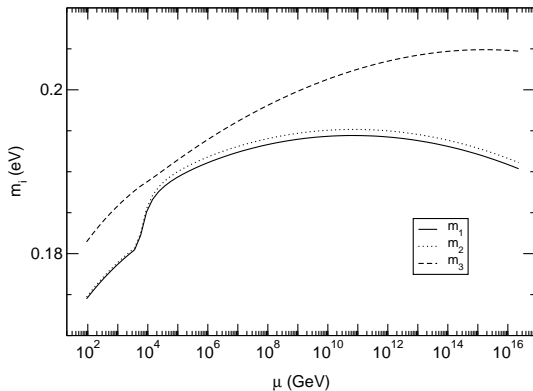
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The RG evolution of neutrino mixing angles θ_{ij}

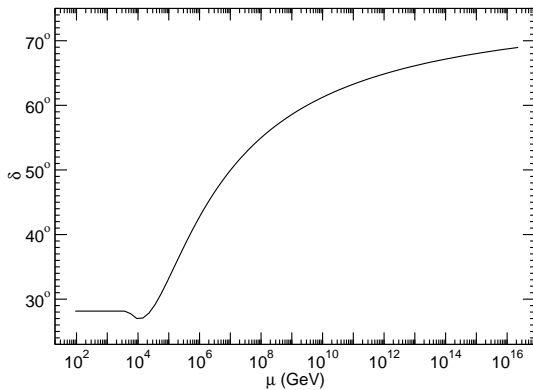


- Quasi-degenerate neutrinos: Large angle magnification in leptonic sector

The RG evolution of neutrino masses m_i



The RG evolution of Dirac Phase



Numerical results

- Bottom-up running:

$$\theta_{12}^{0,q} = 13.02^\circ, \theta_{13}^{0,q} = 0.17^\circ, \theta_{23}^{0,q} = 2.03^\circ \text{ and } \delta_{CP}^{0,q} = 68.93^\circ.$$

- Following HSMU, the neutrino mixing parameters at unification scale are taken to be same as those of quark mixing parameters.

- At unification scale, we choose:

$$m_2^0 = 0.1912 \text{ eV}, \Delta m_{21}^2 = 2.8478 \times 10^{-4} \text{ eV}^2, \\ \Delta m_{32}^2 = 5.3602 \times 10^{-3} \text{ eV}^2.$$

- Top-down running:

$$\theta_{12} = 31.20^\circ, \theta_{13} = 7.22^\circ, \theta_{23} = 50.39^\circ, \delta_{CP} = 28.14^\circ, \\ J_{CP} = 0.102, m_2 = 0.1747 \text{ eV}, \Delta m_{sol}^2 = 7.750 \times 10^{-5} \text{ eV}^2, \\ \Delta m_{atm}^2 = 2.399 \times 10^{-3} \text{ eV}^2.$$

- All low scale parameters are within their $3\text{-}\sigma$ range.
- Threshold corrections are NOT required.
- The mean mass $m = 0.1769 \text{ eV}$ and the “averaged electron neutrino mass” $m_\beta = 0.1747 \text{ eV}$ (slightly below the present reach of KATRIN experiment).

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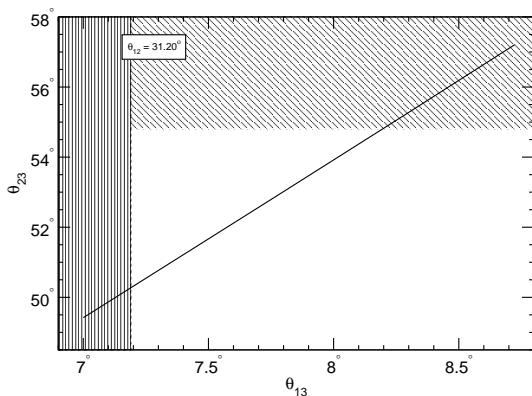
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 $J_{CP} = 0.102$, $m_2 = 0.1747 \text{ eV}$, $\Delta m_{sol}^2 = 7.750 \times 10^{-5} \text{ eV}^2$,
 $\Delta m_{atm}^2 = 2.399 \times 10^{-3} \text{ eV}^2$.
- All low scale parameters are within their $3\text{-}\sigma$ range.
- Threshold corrections are NOT required.
- The mean mass $m = 0.1769 \text{ eV}$ and the “averaged electron neutrino mass” $m_\beta = 0.1747 \text{ eV}$ (slightly below the present reach of KATRIN experiment).

Octant of θ_{23}



- Correlated RG evolution of θ_{13} and θ_{23} : θ_{23} non maximal and in second octant

Summary, so far

- HSMU can be realized for both Dirac as well as Majorana neutrinos
- Several important predictions for Dirac case:
 - Dirac nature: No neutrinoless double beta decay
 - "Averaged electron neutrino mass" m_β : Slightly below KATRIN's proposed sensitivity
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- 1 Introduction
- 2 High Scale Mixing Unification Hypothesis
- 3 Majorana case
- 4 Dirac Case
- 5 Scale of HSMU and SUSY**
- 6 Effect of Phases
- 7 Testing HSMU Hypothesis
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Variation of HSMU and SUSY Breaking Scale

- So far we assumed HSMU to be realized at GUT scale
- HSMU does not depend on “details” of GUT scale theory
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- Similarly SUSY breaking scale and $\tan \beta$ were taken as 5 TeV and 55, respectively
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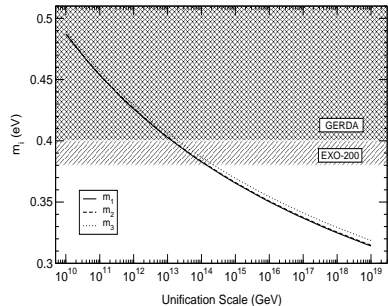
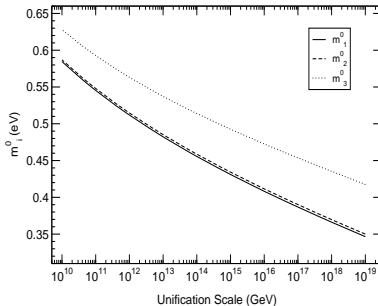
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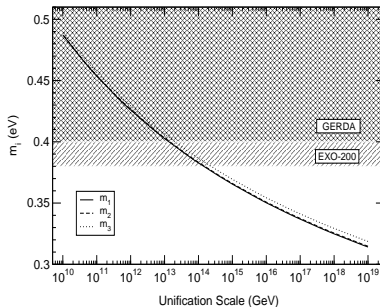
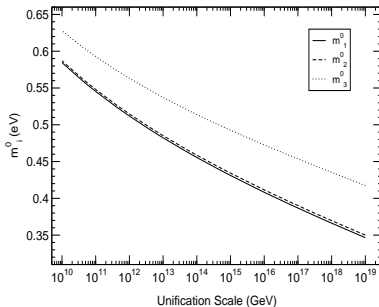
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- Here, we have taken $M_{SUSY} = 5 \text{ TeV}$ and $\tan \beta = 55$.
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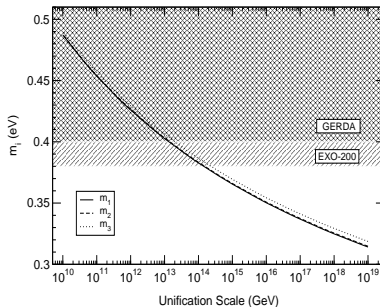
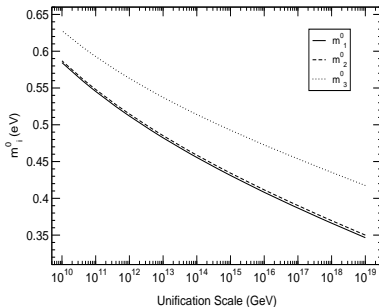
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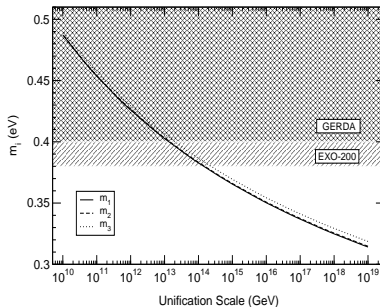
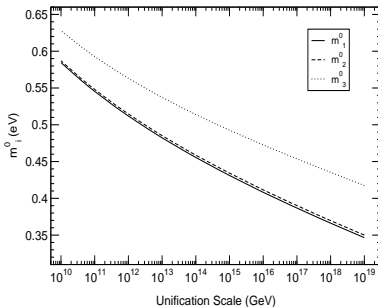
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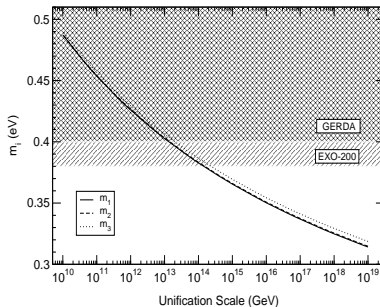
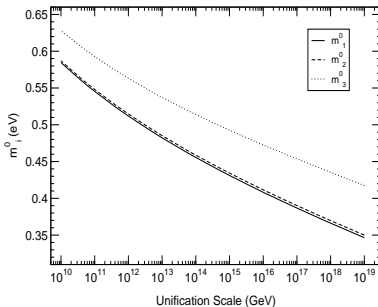
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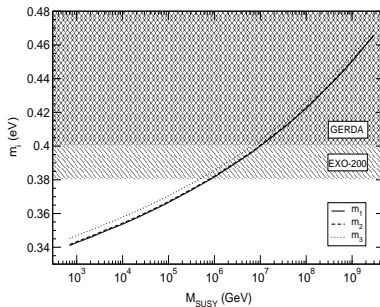
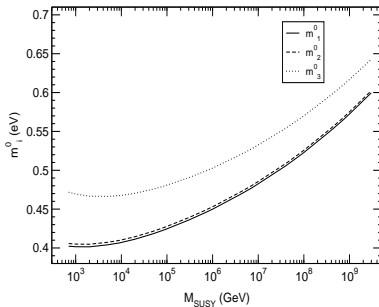
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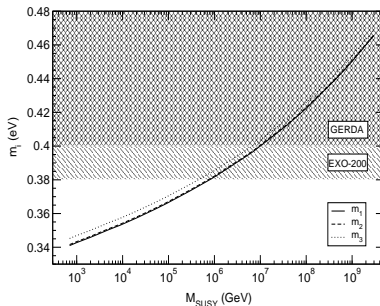
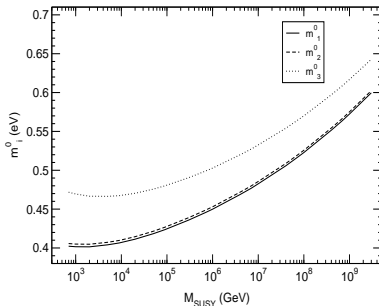
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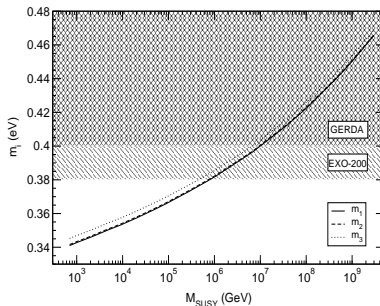
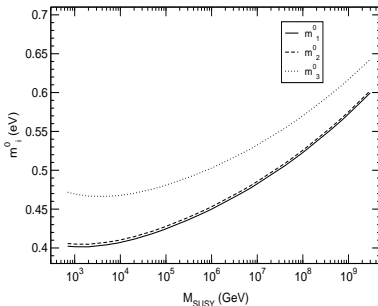
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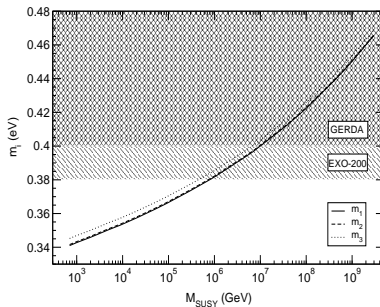
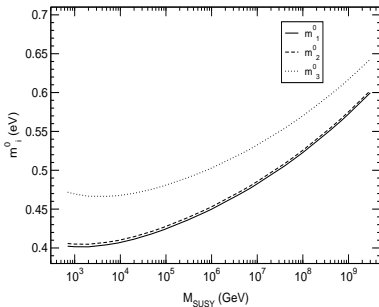
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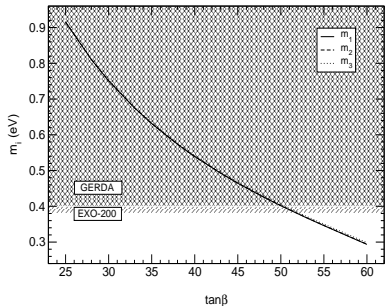
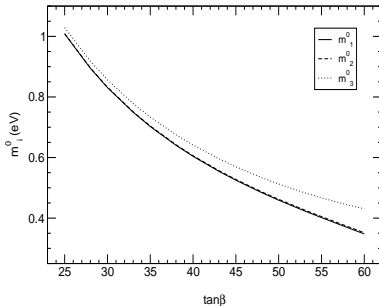
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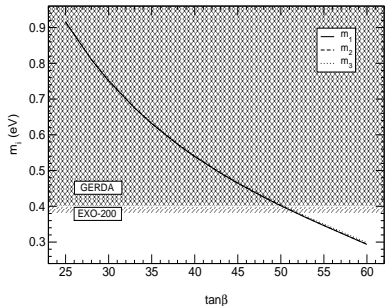
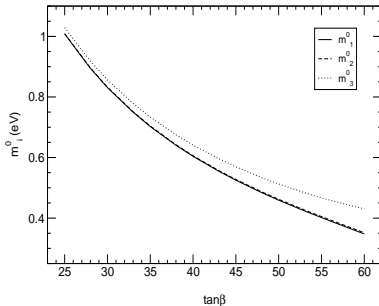
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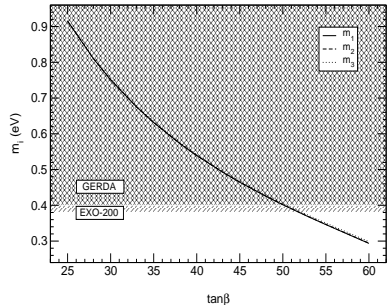
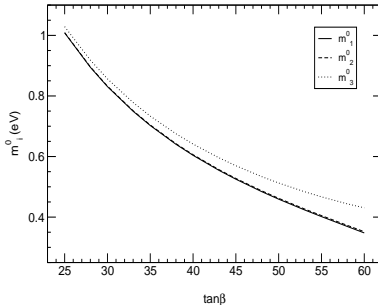
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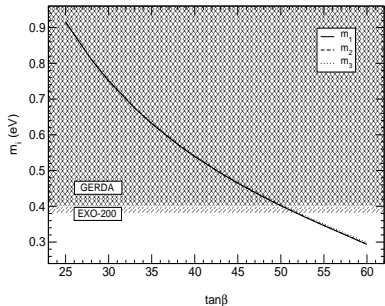
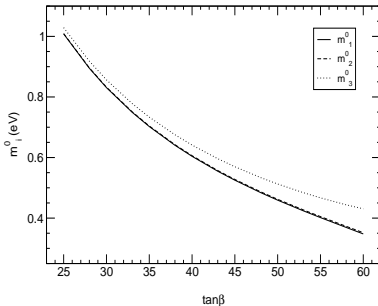
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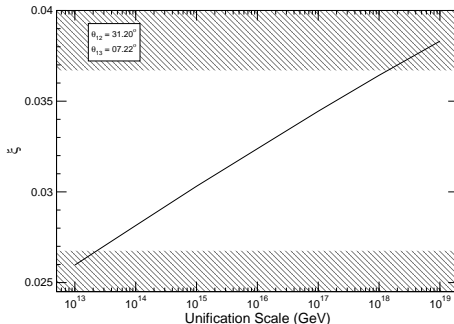
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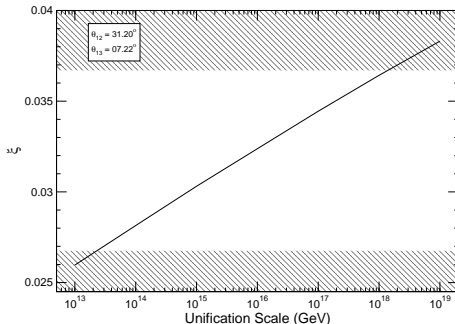
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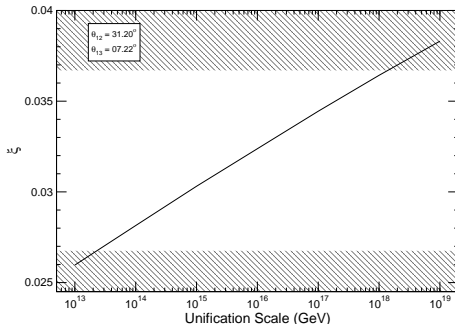
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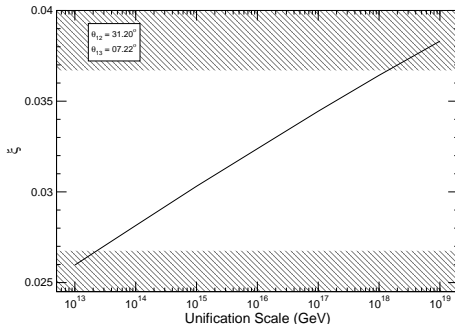
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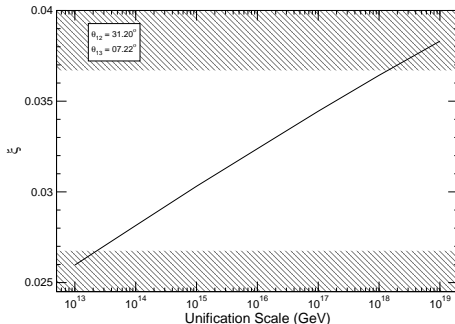
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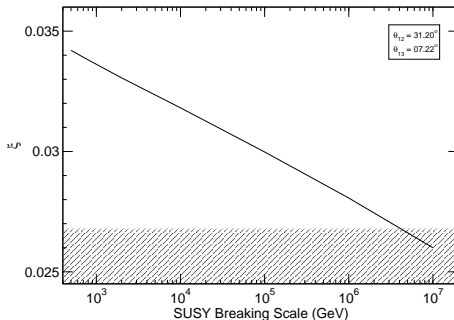
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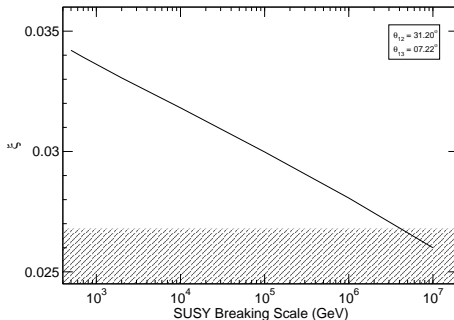
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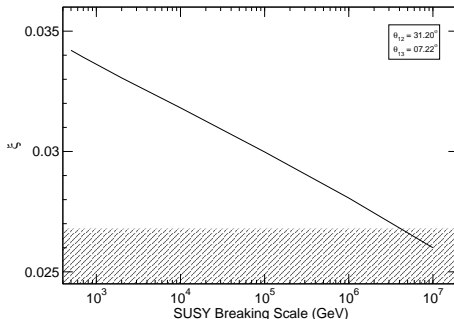
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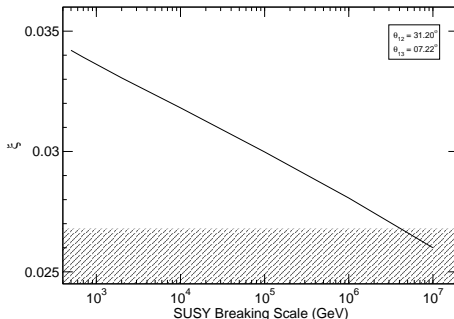
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Dirac Neutrinos: CP Conserving Scenario

- **Plausible Scenario: No CP violation in leptonic sector**
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Dirac Neutrinos: CP Conserving Scenario

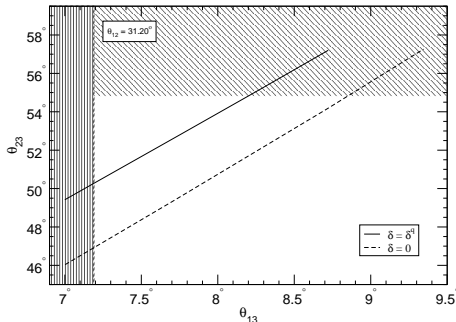
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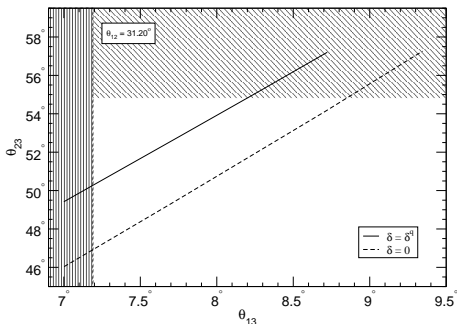
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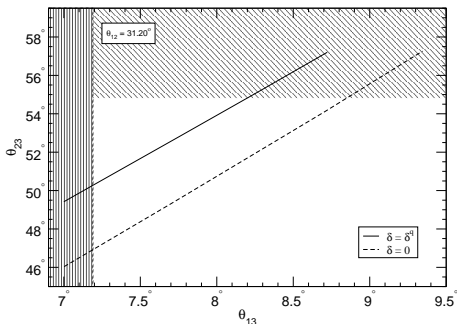
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Majorana Case: CP Violating Scenario

Case I: $\delta_{CP} = \delta_{CP}^q$, $\phi_1 = \phi_2 = 0$

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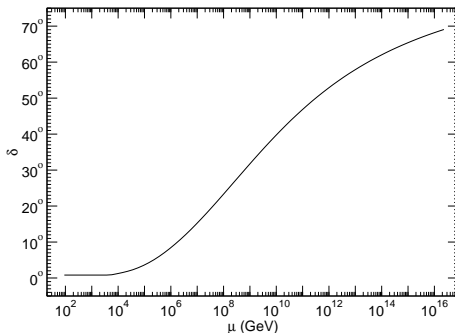
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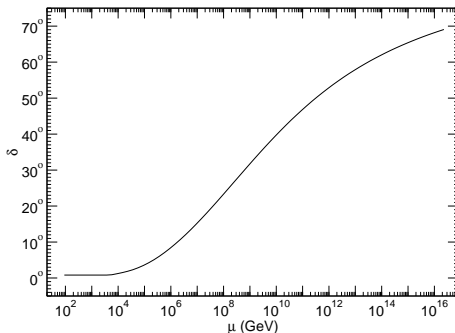
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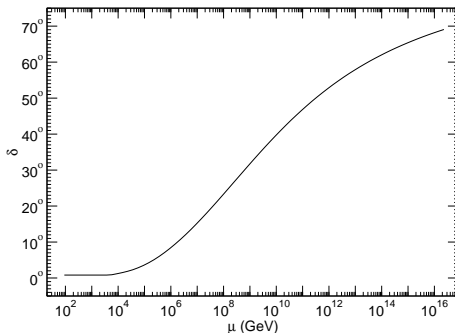
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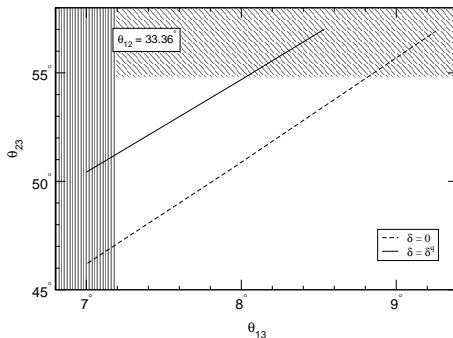
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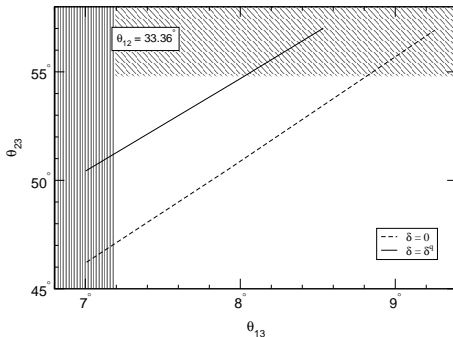
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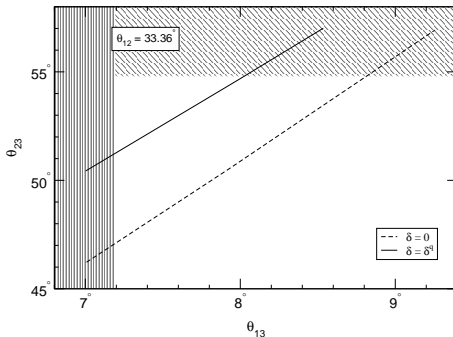
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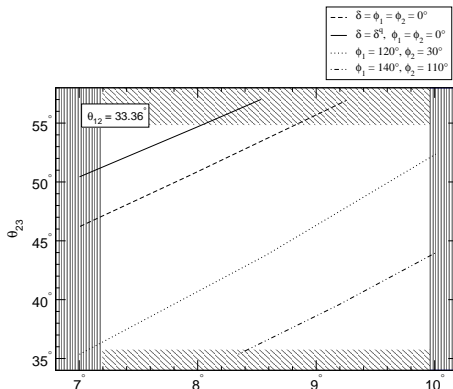
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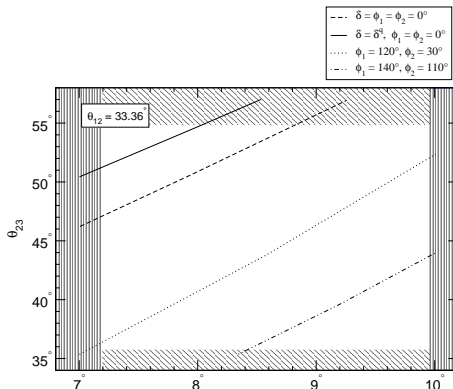
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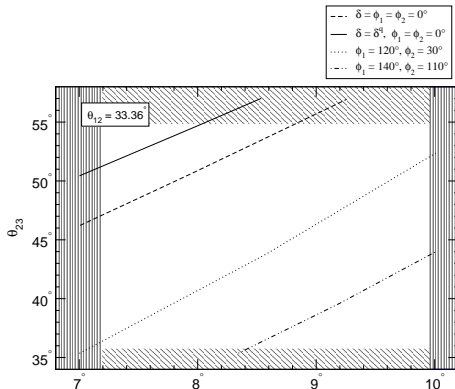
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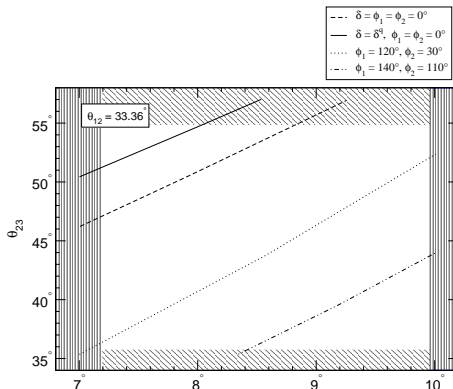
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- 1 Introduction
- 2 High Scale Mixing Unification Hypothesis
- 3 Majorana case
- 4 Dirac Case
- 5 Scale of HSMU and SUSY
- 6 Effect of Phases
- 7 Testing HSMU Hypothesis
- 8 Conclusion and Future Work

Testing HSMU Hypothesis

- HSMU is quite predictive
- Several experiments can test its predictions

Experiment	Majorana	Dirac
$m_{\beta\beta}$ (observed)	✓	×
$m_{\beta\beta} < 0.1$ eV	×	✓
m_{β} (observed) KATRIN	✓	✓
m_{β} (not observed) KATRIN	×	✓
$\theta_{23} > 45^\circ$	✓	✓
$\theta_{23} < 45^\circ$	✓	×
Mass Hierarchy (Normal)	✓	✓
Mass Hierarchy (Inverted)	×	×

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Conclusions and Future Work

- High Scale Mixing Unification (HSMU) of PMNS and CKM parameters is an interesting possibility
- It can be realized with both Dirac and Majorana type neutrinos
- It naturally leads to non zero and “relatively large” values of θ_{13} consistent with present global fits
- It leads to several predictions which can be test by present and near future experiments
- The scale of HSMU is roughly same as that of Grand Unified theories
- This opens up the possibility of realizing HSMU through a GUT
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